

Linear ODEs from an algebraic point of view

Inna Scherbak (joint with Letterio Gatto)

Conference Legacy of Vladimir Arnold

Fields Institute, Toronto

November 24 – 28, 2014

We solve the generic order r linear ODE

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0$$

We solve the generic order r linear ODE

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0$$

e_1, \dots, e_r indeterminate constant coefficients,

$$u(t) \in B_r[[t]], \quad B_r = \mathbb{Q}[e_1, \dots, e_r].$$

We solve the generic order r linear ODE

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0$$

e_1, \dots, e_r indeterminate constant coefficients,

$$u(t) \in B_r[[t]], \quad B_r = \mathbb{Q}[e_1, \dots, e_r].$$

Theorem 1: Let $h_j \in B_r$, $j \in \mathbb{Z}$, be given by

$$\sum_{j \in \mathbb{Z}} h_j z^j = \frac{1}{1 - e_1 z + e_2 z^2 - \dots + (-1)^r e_r z^r}.$$

Then $u_{-j}(t) = \sum_{n \geq j} h_{n-j} \frac{t^n}{n!}$, $0 \leq j \leq r-1$,

is a fundamental system of solutions.

We solve the generic order r linear ODE

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0$$

e_1, \dots, e_r indeterminate constant coefficients,

$$u(t) \in B_r[[t]], \quad B_r = \mathbb{Q}[e_1, \dots, e_r].$$

Theorem 1: Let $h_j \in B_r$, $j \in \mathbb{Z}$, be given by

$$\sum_{j \in \mathbb{Z}} h_j z^j = \frac{1}{1 - e_1 z + e_2 z^2 - \dots + (-1)^r e_r z^r}.$$

$$\text{Then } u_{-j}(t) = \sum_{n \geq j} h_{n-j} \frac{t^n}{n!}, \quad 0 \leq j \leq r-1,$$

is a fundamental system of solutions.

$u_{1-r}(t)$ generates this fundamental system: $u_{-j}(t) = u_{1-r}^{(r-1-j)}(t)$.

We solve the generic order r linear ODE

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0$$

e_1, \dots, e_r indeterminate constant coefficients,

$$u(t) \in B_r[[t]], \quad B_r = \mathbb{Q}[e_1, \dots, e_r].$$

Theorem 1: Let $h_j \in B_r$, $j \in \mathbb{Z}$, be given by

e_1, \dots, e_r the elementary sym functions in ξ_1, \dots, ξ_r



$h_j, j \in \mathbb{Z}$, the complete sym functions: $h_j = 0$ ($j < 0$),
 $h_0 = 1, h_1 = e_1, h_2 = e_1^2 - e_2, \dots$

$$\sum_{j \in \mathbb{Z}} h_j z^j = \frac{1}{1 - e_1 z + e_2 z^2 - \dots + (-1)^r e_r z^r}.$$

$$\text{Then } u_{-j}(t) = \sum_{n \geq j} h_{n-j} \frac{t^n}{n!}, \quad 0 \leq j \leq r-1,$$

is a fundamental system of solutions.

$$u_{1-r}(t) \text{ generates this fundamental system: } u_{-j}(t) = u_{1-r}^{(r-1-j)}(t).$$

$$u^{(r)}(t)-e_1 u^{(r-1)}(t)+e_2 u^{(r-2)}(t)-...+(-1)^r e_r u(t)=0 \Longleftrightarrow X' = M_r X$$

$$x_j(t) = u^{(j)}(t)$$

$$X=\begin{pmatrix}x_0(t)\\ \vdots \\ x_{r-1}(t)\end{pmatrix}$$

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0 \Leftrightarrow X' = M_r X$$

$$M_r = \begin{pmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & 1 \\ (-1)^{r-1} e_r & (-1)^{r-2} e_{r-1} & (-1)^{r-3} e_{r-2} & \dots & e_1 \end{pmatrix} \quad x_j(t) = u^{(j)}(t)$$

$$X = \begin{pmatrix} x_0(t) \\ \vdots \\ x_{r-1}(t) \end{pmatrix}$$

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0 \Leftrightarrow X' = M_r X$$

$$M_r = \begin{pmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & 1 \\ (-1)^{r-1} e_r & (-1)^{r-2} e_{r-1} & (-1)^{r-3} e_{r-2} & \dots & e_1 \end{pmatrix} \quad x_j(t) = u^{(j)}(t)$$

$$X = \begin{pmatrix} x_0(t) \\ \vdots \\ x_{r-1}(t) \end{pmatrix}$$

$$X(t) = \exp(M_r t) \cdot C, \quad C = \begin{pmatrix} c_0 \\ \vdots \\ c_{r-1} \end{pmatrix} = X(0)$$

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0 \Leftrightarrow X' = M_r X$$

$$M_r = \begin{pmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & 1 \\ (-1)^{r-1} e_r & (-1)^{r-2} e_{r-1} & (-1)^{r-3} e_{r-2} & \dots & e_1 \end{pmatrix} \quad x_j(t) = u^{(j)}(t)$$

$$X = \begin{pmatrix} x_0(t) \\ \vdots \\ x_{r-1}(t) \end{pmatrix}$$

$$u(t) = x_0(t), \quad u^{(j)}(0) = c_j \quad \Leftrightarrow \quad X(t) = \exp(M_r t) \cdot C, \quad C = \begin{pmatrix} c_0 \\ \vdots \\ c_{r-1} \end{pmatrix} = X(0)$$

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0 \Leftrightarrow X' = M_r X$$

$$M_r = \begin{pmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & 1 \\ (-1)^{r-1} e_r & (-1)^{r-2} e_{r-1} & (-1)^{r-3} e_{r-2} & \dots & e_1 \end{pmatrix} \quad x_j(t) = u^{(j)}(t)$$

$$X = \begin{pmatrix} x_0(t) \\ \vdots \\ x_{r-1}(t) \end{pmatrix}$$

$$u(t) = x_0(t), \quad u^{(j)}(0) = c_j \quad \Leftrightarrow \quad X(t) = \exp(M_r t) \cdot C, \quad C = \begin{pmatrix} c_0 \\ \vdots \\ c_{r-1} \end{pmatrix} = X(0)$$

Remark: $\exp(M_r t)$ is the Wronski matrix of the standard

fundamental system of the solutions to the ODE, $v_0(t), \dots, v_{r-1}(t)$, that is $v_i^{(j)}(0) = \delta_{ij}$, $0 \leq i, j \leq r-1$ (standard initial conditions).

Theorem 2: The last column of $\exp(M_r t)$ is our universal fundamental system:

$$\exp(M_r t) = \left(\begin{array}{ccc|c} v_0 & v_1 & \cdots & v_{r-1} = u_{1-r} \\ \vdots & \vdots & \ddots & v_{r-1} = u_{2-r} \\ v_0^{(r-1)} & v_1^{(r-1)} & \cdots & v_{r-1}^{(r-1)} = u_0 \end{array} \right)$$

$$u_0(t) = \sum_{k \geq 0} h_k \frac{t^k}{k!}, \dots, u_{1-r}(t) = \sum_{k \geq r-1} h_{k+1-r} \frac{t^k}{k!}$$

The universal and the standard fundamental systems are related as follows,

$$\begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix} = \begin{pmatrix} 1 & h_1 & h_2 & \cdots & \cdots & h_r \\ 0 & 1 & h_1 & \cdots & \cdots & h_{r-1} \\ 0 & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_r \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_r \end{pmatrix} = \begin{pmatrix} 1 & -e_1 & e_2 & \cdots & \cdots & (-1)^r e_r \\ 0 & 1 & -e_1 & \cdots & \cdots & (-1)^{r-1} e_{r-1} \\ \vdots & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix}$$

The universal and the standard fundamental systems are related as follows,

$$\begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix} = \begin{pmatrix} 1 & h_1 & h_2 & \cdots & \cdots & h_r \\ 0 & 1 & h_1 & \cdots & \cdots & h_{r-1} \\ 0 & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_r \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_r \end{pmatrix} = \begin{pmatrix} 1 & -e_1 & e_2 & \cdots & \cdots & (-1)^r e_r \\ 0 & 1 & -e_1 & \cdots & \cdots & (-1)^{r-1} e_{r-1} \\ \vdots & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix}$$

$$\mathbf{H}\mathbf{E}=\mathbf{I}$$

The universal and the standard fundamental systems are related as follows,

$$\begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix} = \begin{pmatrix} 1 & h_1 & h_2 & \cdots & \cdots & h_r \\ 0 & 1 & h_1 & \cdots & \cdots & h_{r-1} \\ 0 & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_r \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_r \end{pmatrix} = \begin{pmatrix} 1 & -e_1 & e_2 & \cdots & \cdots & (-1)^r e_r \\ 0 & 1 & -e_1 & \cdots & \cdots & (-1)^{r-1} e_{r-1} \\ \vdots & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix}$$

$$\mathbf{H}\mathbf{E}=\mathbf{I}$$

Generating functions of the complete and the elementary symmetric functions,

$$H(z) = \sum_{k \geq 0} h_k z^k = \prod_{i \geq 1} (1 - \xi_i z)^{-1} \quad \text{and} \quad E(z) = \sum_{k=0}^r e_k z^k = \prod_{i=1}^r (1 + \xi_i z)$$

The universal and the standard fundamental systems are related as follows,

$$\begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix} = \begin{pmatrix} 1 & h_1 & h_2 & \cdots & \cdots & h_r \\ 0 & 1 & h_1 & \cdots & \cdots & h_{r-1} \\ 0 & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_r \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_r \end{pmatrix} = \begin{pmatrix} 1 & -e_1 & e_2 & \cdots & \cdots & (-1)^r e_r \\ 0 & 1 & -e_1 & \cdots & \cdots & (-1)^{r-1} e_{r-1} \\ \vdots & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix}$$

$$HE=I$$

Generating functions of the complete and the elementary symmetric functions,

$H(z) = \sum_{k \geq 0} h_k z^k = \prod_{i \geq 1} (1 - \xi_i z)^{-1}$ and $E(z) = \sum_{k=0}^r e_k z^k = \prod_{i=1}^r (1 + \xi_i z)$, satisfy

$$H(-z)E(z)=1$$

Corollary: The (unique) solution to the *Cauchy problem* ,

$$u^{(r)} - e_1 u^{(r-1)} + e_2 u^{(r-2)} - \dots + (-1)^r e_r u = 0, \quad u^{(j)}(0) = c_j \in B_r,$$

$$u_C(t) = \begin{pmatrix} c_0 & \cdots & c_{r-1} \end{pmatrix} \begin{pmatrix} 1 & -e_1 & e_2 & \cdots & \cdots & (-1)^r e_r \\ 0 & 1 & -e_1 & \cdots & \cdots & (-1)^{r-1} e_{r-1} \\ \vdots & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} \textcolor{magenta}{u}_0 \\ \textcolor{magenta}{u}_{-1} \\ \vdots \\ \vdots \\ \textcolor{magenta}{u}_{1-r} \end{pmatrix}.$$

Corollary: The (unique) solution to the *Cauchy problem*,

$$u^{(r)} - e_1 u^{(r-1)} + e_2 u^{(r-2)} - \dots + (-1)^r e_r u = 0, \quad u^{(j)}(0) = c_j \in B_r,$$

$$u_C(t) = \begin{pmatrix} c_0 & \cdots & c_{r-1} \end{pmatrix} \begin{pmatrix} 1 & -e_1 & e_2 & \cdots & \cdots & (-1)^r e_r \\ 0 & 1 & -e_1 & \cdots & \cdots & (-1)^{r-1} e_{r-1} \\ \vdots & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix}.$$

Solution in the non-homogeneous case, $g(t) = \sum_{k \geq 0} b_k \frac{t^k}{k!} \in B_r[[t]]$:

$u_C(t) + p(t)$, where $p(t)$ the particular solution with vanishing initial conditions.

Corollary: The (unique) solution to the *Cauchy problem* ,

$$u^{(r)} - e_1 u^{(r-1)} + e_2 u^{(r-2)} - \dots + (-1)^r e_r u = 0, \quad u^{(j)}(0) = c_j \in B_r,$$

$$u_C(t) = \begin{pmatrix} c_0 & \cdots & c_{r-1} \end{pmatrix} \begin{pmatrix} 1 & -e_1 & e_2 & \cdots & \cdots & (-1)^r e_r \\ 0 & 1 & -e_1 & \cdots & \cdots & (-1)^{r-1} e_{r-1} \\ \vdots & 0 & 1 & \cdots & \cdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \cdots & 1 \end{pmatrix} \begin{pmatrix} u_0 \\ u_{-1} \\ \vdots \\ u_{1-r} \end{pmatrix}.$$

Solution in the non-homogeneous case, $g(t) = \sum_{k \geq 0} b_k \frac{t^k}{k!} \in B_r[[t]]$:

$u_C(t) + p(t)$, where $p(t)$ the particular solution with vanishing initial conditions.

Theorem 3: We have $p(t) = \sum_{k \geq r} p_k \frac{t^k}{k!}$, where $p_k \in B_r$, $k \geq r$, are given by

$$\sum_{k \geq r} p_k z^{k-r} = \frac{\sum_{k \geq 0} b_k z^k}{1 - e_1 z + e_2 z^2 - \dots + (-1)^r e_r z^r}.$$

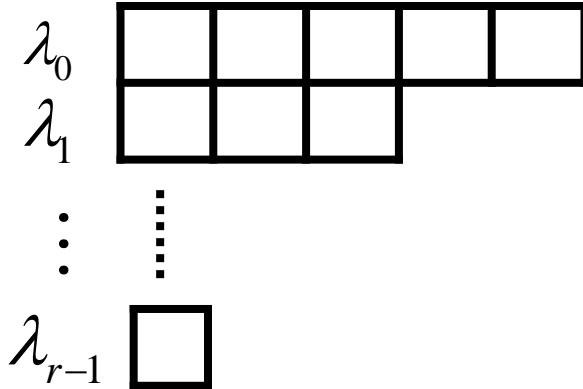
Schur function:

$$S_\lambda(\xi) = \det(\xi_{\lambda_j+i-j})_{0 \leq i, j \leq r-1},$$

where

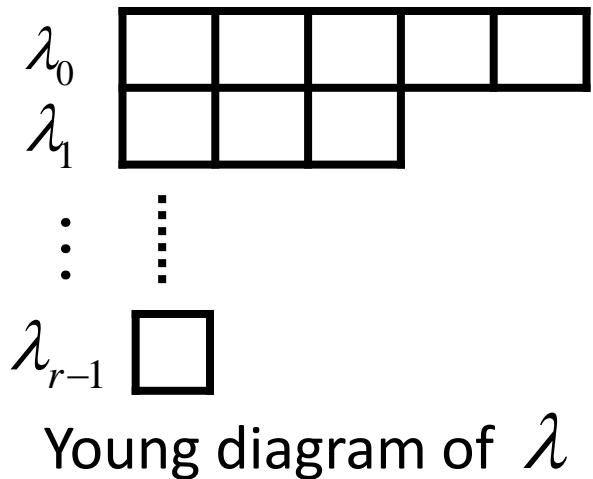
$$\lambda = (\lambda_0 \geq \lambda_1 \geq \dots \geq \lambda_{r-1} \geq 0) \in P_r,$$

a partition of $|\lambda| = \sum \lambda_j$ of length $\leq r$,
 $\{\xi_i\}_{i \in \mathbb{Z}}$ variables.



Young diagram of λ

Schur function:



$$S_\lambda(\xi) = \det(\xi_{\lambda_j+i-j})_{0 \leq i,j \leq r-1},$$

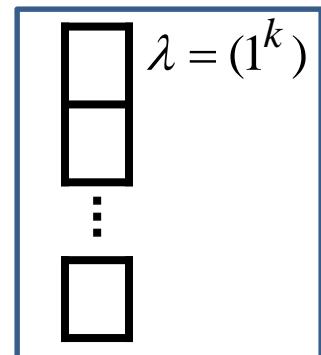
where

$$\lambda = (\lambda_0 \geq \lambda_1 \geq \dots \geq \lambda_{r-1} \geq 0) \in P_r,$$

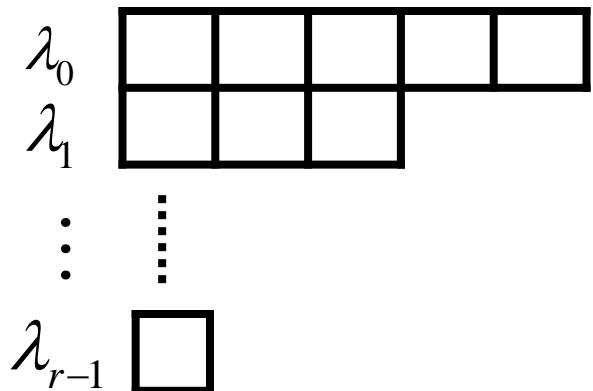
a partition of $|\lambda| = \sum \lambda_j$ of length $\leq r$,
 $\{\xi_i\}_{i \in \mathbb{Z}}$ variables.

Examples: The coefficients e_1, \dots, e_r of the ODE and the coefficients $h = \{h_j\}_{j \in \mathbb{Z}}$ of the universal solution are related by the Giambelli formula

$$e_k = S_{(1^k)}(h).$$



Schur function:



$$S_\lambda(\xi) = \det(\xi_{\lambda_j+i-j})_{0 \leq i,j \leq r-1},$$

where

$$\lambda = (\lambda_0 \geq \lambda_1 \geq \dots \geq \lambda_{r-1} \geq 0) \in P_r,$$

a partition of $|\lambda| = \sum \lambda_j$ of length $\leq r$,
 $\{\xi_i\}_{i \in \mathbb{Z}}$ variables.

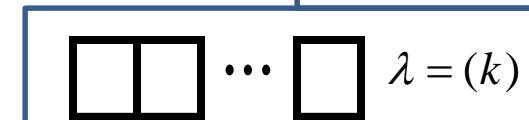
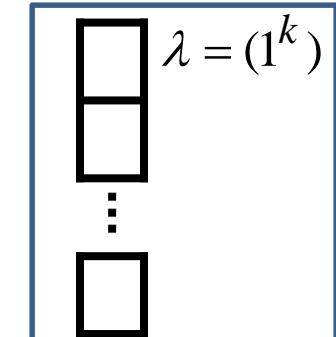
Examples: The coefficients e_1, \dots, e_r of the ODE and the coefficients $h = \{h_j\}_{j \in \mathbb{Z}}$ of the universal solution are related by the Giambelli formula

$$e_k = S_{(1^k)}(h).$$

As the elementary and the complete symmetric functions

in $\xi_1, \xi_2, \dots, \xi_r$, $e_k = S_{(1^k)}(\xi)$, $h_k = S_{(k)}(\xi)$, where

$$\xi = \{\xi_j\}_{j \in \mathbb{Z}}, \quad \xi_0 = 1, \quad \xi_j = 0, \quad j < 0, \quad j > r.$$



For $f_j = f_j(t) \in B_r[[t]]$ ($0 \leq j \leq r-1$), denote $\bar{f} = (f_0, f_1, \dots, f_{r-1})$.

The Wronskian of \bar{f} : $W[\bar{f}](t) = \det(f_j^{(i)})_{0 \leq i, j \leq r-1}$.

For $f_j = f_j(t) \in B_r[[t]]$ ($0 \leq j \leq r-1$), denote $\bar{f} = (f_0, f_1, \dots, f_{r-1})$.

The Wronskian of \bar{f} : $W[\bar{f}](t) = \det(f_j^{(i)})_{0 \leq i, j \leq r-1}$.

Motivation: The role Wronskians play in Schubert calculus on Grassmannian
[papers of L. Goldberg; A. Eremenko and A. Gabrielov; B. and M. Shapiro;
E. Mukhin, V. Tarasov, and A. Varchenko; ...]

For $f_j = f_j(t) \in B_r[[t]]$ ($0 \leq j \leq r-1$), denote $\bar{f} = (f_0, f_1, \dots, f_{r-1})$.

The Wronskian of \bar{f} : $W[\bar{f}](t) = \det(f_j^{(i)})_{0 \leq i, j \leq r-1}$.

Motivation: The role Wronskians play in Schubert calculus on Grassmannian [papers of L. Goldberg; A. Eremenko and A. Gabrielov; B. and M. Shapiro; E. Mukhin, V. Tarasov, and A. Varchenko; ...]

Generalized Wronskian of \bar{f} corresponding to $\lambda = (\lambda_0 \geq \lambda_1 \geq \dots \geq \lambda_{r-1} \geq 0)$:

$$W[\lambda, \bar{f}](t) = \det(f_j^{(\lambda_{r-1-i} + i)})_{0 \leq i, j \leq r-1}.$$

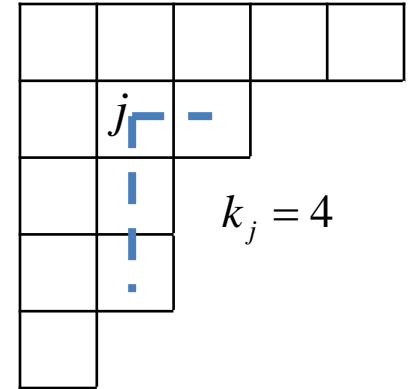
Theorem 4 (Derivatives of the Wronskian in terms of the generalized Wronskians):

$$W[\bar{f}]^{(k)}(t) = \sum_{|\lambda|=k} c_\lambda W[\lambda, \bar{f}](t), \quad \text{where } c_\lambda = \frac{|\lambda|!}{k_1 \cdots k_{|\lambda|}}.$$

Theorem 4 (Derivatives of the Wronskian in terms of the generalized Wronskians):

$$W[\bar{f}]^{(k)}(t) = \sum_{|\lambda|=k} c_\lambda W[\lambda, \bar{f}](t), \text{ where } c_\lambda = \frac{|\lambda|!}{k_1 \cdots k_{|\lambda|}}.$$

(hook formula)



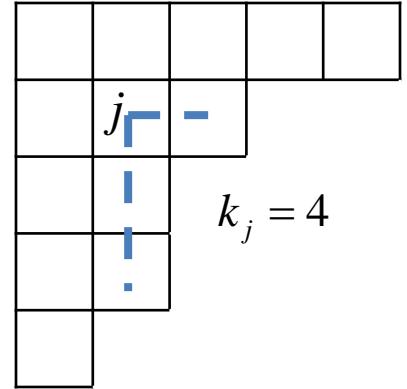
Theorem 4 (Derivatives of the Wronskian in terms of the generalized Wronskians):

$$W[\bar{f}]^{(k)}(t) = \sum_{|\lambda|=k} c_\lambda W[\lambda, \bar{f}](t), \quad \text{where } c_\lambda = \frac{|\lambda|!}{k_1 \cdots k_{|\lambda|}}.$$

(hook formula)

Let now $\bar{f} = (f_0, f_1, \dots, f_{r-1})$ be
a fundamental system of the ODE

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0.$$



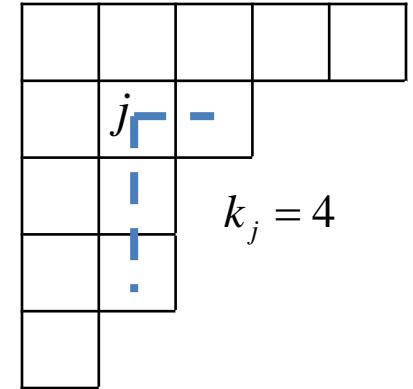
Theorem 4 (Derivatives of the Wronskian in terms of the generalized Wronskians):

$$W[\bar{f}]^{(k)}(t) = \sum_{|\lambda|=k} c_\lambda W[\lambda, \bar{f}](t), \quad \text{where } c_\lambda = \frac{|\lambda|!}{k_1 \cdots k_{|\lambda|}}.$$

(hook formula)

Let now $\bar{f} = (f_0, f_1, \dots, f_{r-1})$ be
a fundamental system of the ODE

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0.$$



Theorem 5: We have $W[\lambda, \bar{f}] = S_\lambda(h) W[\bar{f}]$ [Giambelli's formula],

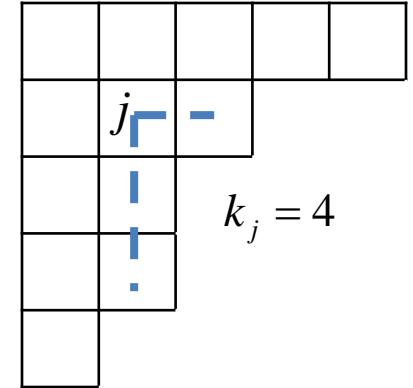
Theorem 4 (Derivatives of the Wronskian in terms of the generalized Wronskians):

$$W[\bar{f}]^{(k)}(t) = \sum_{|\lambda|=k} c_\lambda W[\lambda, \bar{f}](t), \quad \text{where } c_\lambda = \frac{|\lambda|!}{k_1 \cdots k_{|\lambda|}}.$$

(hook formula)

Let now $\bar{f} = (f_0, f_1, \dots, f_{r-1})$ be
a fundamental system of the ODE

$$u^{(r)}(t) - e_1 u^{(r-1)}(t) + e_2 u^{(r-2)}(t) - \dots + (-1)^r e_r u(t) = 0.$$



Theorem 5: We have $W[\lambda, \bar{f}] = S_\lambda(h) W[\bar{f}]$ [Giambelli's formula],

$$h_k W[\lambda, \bar{f}] = \sum_{\mu} W[\mu, \bar{f}] \quad [\text{Pieri's formula}],$$

the sum over all partitions $\mu \in P_r$ satisfying

$$|\mu| = k + |\lambda|, \quad \mu_0 \geq \lambda_0 \geq \mu_1 \geq \lambda_1 \geq \dots \geq \mu_{r-1} \geq \lambda_{r-1}.$$

Denote $\underline{\mathcal{W}}_r$ the free \mathbb{Z} -module generated by $\{W[\lambda, \bar{f}], \lambda \in P_r\}$,
where \bar{f} a fundamental system of the ODE.

Denote $\underline{\mathcal{W}}_r$ the free \mathbb{Z} -module generated by $\{W[\lambda, \bar{f}], \lambda \in P_r\}$,
where \bar{f} a fundamental system of the ODE.

Let $G = G(r, \mathbb{P}^\infty)$ be the Grassmannian, $\sigma_\lambda \in H^*(G, \mathbb{Z})$ Schubert classes,
 $\Omega_\lambda \in H_*(G, \mathbb{Z})$ the Poincare duals.

Denote $\underline{\mathcal{W}}_r$ the free \mathbb{Z} -module generated by $\{W[\lambda, \bar{f}], \lambda \in P_r\}$,
where \bar{f} a fundamental system of the ODE.

Let $G = G(r, \mathbb{P}^\infty)$ be the Grassmannian, $\sigma_\lambda \in H^*(G, \mathbb{Z})$ Schubert classes,
 $\Omega_\lambda \in H_*(G, \mathbb{Z})$ the Poincare duals.

Mapping $h_j \mapsto \sigma_{(j)}$ defines a surjection $B_r \rightarrow H^*(G, \mathbb{Z})$, and its kernel gives relations on generators ($B_r = \mathbb{Q}[e_1, \dots, e_r]$ the ring of symmetric functions).

Denote $\underline{\mathcal{W}}_r$ the free \mathbb{Z} -module generated by $\{W[\lambda, \bar{f}], \lambda \in P_r\}$,
 were \bar{f} a fundamental system of the ODE.

Let $G = G(r, \mathbb{P}^\infty)$ be the Grassmannian, $\sigma_\lambda \in H^*(G, \mathbb{Z})$ Schubert classes,
 $\Omega_\lambda \in H_*(G, \mathbb{Z})$ the Poincare duals.

Mapping $h_j \mapsto \sigma_{(j)}$ defines a surjection $B_r \rightarrow H^*(G, \mathbb{Z})$, and its kernel gives relations on generators ($B_r = \mathbb{Q}[e_1, \dots, e_r]$ the ring of symmetric functions).

Well-known Giambelli's and Pieri's formulae of Schubert calculus:

$$\sigma_\lambda = S_\lambda(\bar{\sigma}), \quad \sigma_{(k)} \sigma_\lambda = \sum_\mu \sigma_\mu,$$

where $\bar{\sigma} = (1, \sigma_{(1)}, \sigma_{(2)}, \dots)$, $\sigma_{(j)}$ special Schubert classes.

Denote $\underline{\mathcal{W}}_r$ the free \mathbb{Z} -module generated by $\{W[\lambda, \bar{f}], \lambda \in P_r\}$,
 were \bar{f} a fundamental system of the ODE.

Let $G = G(r, \mathbb{P}^\infty)$ be the Grassmannian, $\sigma_\lambda \in H^*(G, \mathbb{Z})$ Schubert classes,
 $\Omega_\lambda \in H_*(G, \mathbb{Z})$ the Poincare duals.

Mapping $h_j \mapsto \sigma_{(j)}$ defines a surjection $B_r \rightarrow H^*(G, \mathbb{Z})$, and its kernel gives relations on generators ($B_r = \mathbb{Q}[e_1, \dots, e_r]$ the ring of symmetric functions).

Well-known Giambelli's and Pieri's formulae of Schubert calculus:

$$\sigma_\lambda = S_\lambda(\bar{\sigma}), \quad \sigma_{(k)} \sigma_\lambda = \sum_\mu \sigma_\mu,$$

where $\bar{\sigma} = (1, \sigma_{(1)}, \sigma_{(2)}, \dots)$, $\sigma_{(j)}$ special Schubert classes.

Corollary: The \mathbb{Z} -module isomorphism $\underline{\mathcal{W}}_r \rightarrow H_*(G, \mathbb{Z})$ defined by $W[\lambda, \bar{f}] \mapsto \Omega_\lambda$
 is an isomorphism of $H^*(G, \mathbb{Z})$ -modules.

From B_r and \mathcal{W}_r to *bosonic and fermionic spaces* :

From B_r and \mathcal{W}_r to *bosonic and fermionic spaces* :

We write $u_j(t) = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!}$, $j \in \mathbb{Z}$;

From B_r and \mathcal{W}_r to *bosonic and fermionic spaces* :

We write $u_j(t) = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!}$, $j \in \mathbb{Z}$;

$\bar{u} = (u_0, u_{-1}, \dots, u_{1-r})$ - the universal fundamental system of the ODE;

$K_r = \text{Span}_{\mathbb{Q}}\{u_j(t), j \geq -r+1\} = \text{Span}_{B_r}\{u_0, u_{-1}, \dots, u_{1-r}\} \subseteq B_r[[t]]$

- the sub-module of all the solutions of the ODE;

From B_r and \mathcal{W}_r to *bosonic and fermionic spaces* :

We write $u_j(t) = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!}$, $j \in \mathbb{Z}$;

$\bar{u} = (u_0, u_{-1}, \dots, u_{1-r})$ - the universal fundamental system of the ODE;

$K_r = \text{Span}_{\mathbb{Q}}\{u_j(t), j \geq -r+1\} = \text{Span}_{B_r}\{u_0, u_{-1}, \dots, u_{1-r}\} \subseteq B_r[[t]]$

- the sub-module of all the solutions of the ODE;

Theorem 6: We have

$$\wedge^r K_r = \text{Span}_{B_r}\{\bar{u}_\lambda := u_{\lambda_0} \wedge u_{\lambda_1-1} \wedge \dots \wedge u_{\lambda_{r-1}-r+1}, \lambda \in P_r\};$$

$$\bar{u}_\lambda = S_\lambda(h) \bar{u}_O, \text{ where } \bar{u}_O = u_0 \wedge u_{-1} \wedge \dots \wedge u_{1-r}.$$

From B_r and \mathcal{W}_r to *bosonic and fermionic spaces* :

We write $u_j(t) = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!}$, $j \in \mathbb{Z}$;

$\bar{u} = (u_0, u_{-1}, \dots, u_{1-r})$ - the universal fundamental system of the ODE;

$K_r = \text{Span}_{\mathbb{Q}}\{u_j(t), j \geq -r+1\} = \text{Span}_{B_r}\{u_0, u_{-1}, \dots, u_{1-r}\} \subseteq B_r[[t]]$

- the sub-module of all the solutions of the ODE;

Theorem 6: We have

$$\wedge^r K_r = \text{Span}_{B_r}\{\bar{u}_\lambda := u_{\lambda_0} \wedge u_{\lambda_1-1} \wedge \dots \wedge u_{\lambda_{r-1}-r+1}, \lambda \in P_r\};$$

$$\bar{u}_\lambda = S_\lambda(h)\bar{u}_O, \text{ where } \bar{u}_O = u_0 \wedge u_{-1} \wedge \dots \wedge u_{1-r}.$$

Corollary: The correspondence $\bar{u}_\lambda \leftrightarrow W[\lambda, \bar{u}]$ gives the B_r -module isomorphism $\mathcal{W}_r \sim \wedge^r K_r$.

We call

$B_r = \mathbb{Q}[e_1, \dots, e_r]$ the *r-th Bosonic space*;

$F_0^r := \text{Span}_{B_r} \{ \bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots, \lambda \in P_r \}$ the *r-th Fermionic space*

(of zero total charge)

We call

$B_r = \mathbb{Q}[e_1, \dots, e_r]$ the *r-th Bosonic space*;

$F_0^r := \text{Span}_{B_r} \{ \bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots, \lambda \in P_r \}$ the *r-th Fermionic space*

(of zero total charge)

We have $F_0^r \sim \wedge^r K_r \sim \mathcal{W}_r$.

We call

$B_r = \mathbb{Q}[e_1, \dots, e_r]$ the *r-th Bosonic space*;

$F_0^r := \text{Span}_{B_r} \{ \bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots, \lambda \in P_r \}$ the *r-th Fermionic space*

(of zero total charge)

We have $F_0^r \sim \wedge^r K_r \sim \mathcal{W}_r$.

Remark: F_0^r is well-defined for $r = \infty$ whereas \mathcal{W}_r does not.

We call

$B_r = \mathbb{Q}[e_1, \dots, e_r]$ the *r-th Bosonic space*;

$F_0^r := \text{Span}_{B_r} \{ \bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots, \lambda \in P_r \}$ the *r-th Fermionic space*
(of zero total charge).

We have $F_0^r \sim \wedge^r K_r \sim \mathcal{W}_r$.

Remark: F_0^r is well-defined for $r = \infty$ whereas \mathcal{W}_r does not.

If $r = \infty$ we have $B_\infty = \mathbb{Q}[e_1, e_2, \dots]$, $K_\infty = \text{Span}_{B_\infty} \{u_0, u_{-1}, \dots\}$;

$u_j(t) = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!}$, $j \in \mathbb{Z}$, all are solutions to the “*ODE of infinite order*”,

where $\{h_i\}_{i \in \mathbb{Z}}$ are defined in the same way as for $r < \infty$,

$$E_\infty(t) \sum_{n \in \mathbb{Z}} h_n t^n = 1, \quad E_\infty(t) = 1 + \sum_{j \geq 1} (-1)^j e_j t^j.$$

The Boson-Fermion correspondence

The Boson-Fermion correspondence

is the B_r -module isomorphism $c_0^r : F_0^r \rightarrow B_r$ defined by

$$\bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots \mapsto S_\lambda(h)$$

[this terminology is standard if $r = \infty$, see V. Kac, A.K. Raina].

The Boson-Fermion correspondence

is the B_r -module isomorphism $c_0^r : F_0^r \rightarrow B_r$ defined by

$$\bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots \mapsto S_\lambda(h)$$

[this terminology is standard if $r = \infty$, see V. Kac, A.K. Raina] .

Denote $\{x_j\}_{j \geq 1}$ the elements of B_r defined by the generating function

$$\sum_{j \geq 1} x_j z^j = -\log(1 - e_1 z + e_2 z^2 - \dots + (-1)^r e_r z^r).$$

In terms of symmetric functions, $jx_j = \xi_1^j + \dots + \xi_r^j$.

The Boson-Fermion correspondence

is the B_r -module isomorphism $c_0^r : F_0^r \rightarrow B_r$ defined by

$$\bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots \mapsto S_\lambda(h)$$

[this terminology is standard if $r = \infty$, see V. Kac, A.K. Raina].

Denote $\{x_j\}_{j \geq 1}$ the elements of B_r defined by the generating function

$$\sum_{j \geq 1} x_j z^j = -\log(1 - e_1 z + e_2 z^2 - \dots + (-1)^r e_r z^r).$$

In terms of symmetric functions, $jx_j = \xi_1^j + \dots + \xi_r^j$.

We have $B_r = \mathbb{Q}[x_1, \dots, x_r]$.

Thus e_1, \dots, e_r, h_i ($i \geq 1$) all are polynomials in x_1, \dots, x_r .

The Boson-Fermion correspondence

is the B_r -module isomorphism $c_0^r : F_0^r \rightarrow B_r$ defined by

$$\bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots \mapsto S_\lambda(h)$$

[this terminology is standard if $r = \infty$, see V. Kac, A.K. Raina].

Denote $\{x_j\}_{j \geq 1}$ the elements of B_r defined by the generating function

$$\sum_{j \geq 1} x_j z^j = -\log(1 - e_1 z + e_2 z^2 - \dots + (-1)^r e_r z^r).$$

In terms of symmetric functions, $jx_j = \xi_1^j + \dots + \xi_r^j$.

We have $B_r = \mathbb{Q}[x_1, \dots, x_r]$.

Thus e_1, \dots, e_r, h_i ($i \geq 1$) all are polynomials in x_1, \dots, x_r .

If $r = \infty$, then $\sum_{n \geq 1} x_n t^n = -\log E_\infty(t)$, $E_\infty(t) = 1 + \sum_{j \geq 1} (-1)^j e_j t^j$.

We write now $u_j = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!} = u_j(x_1, \dots, x_r; t)$, $j \in \mathbb{Z}$.

We write now $u_j = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!} = u_j(x_1, \dots, x_r; t)$, $j \in \mathbb{Z}$.

Operator $\partial_i = \partial / \partial x_i : B_r \rightarrow B_r$ induces a \mathbb{Q} -linear map

$$\partial_i : K_r \rightarrow K_r \quad (i = 1, \dots, r).$$

We write now $u_j = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!} = u_j(x_1, \dots, x_r; t)$, $j \in \mathbb{Z}$.

Operator $\partial_i = \partial / \partial x_i : B_r \rightarrow B_r$ induces a \mathbb{Q} -linear map

$$\partial_i : K_r \rightarrow K_r \quad (i = 1, \dots, r).$$

Operator $D^k = \frac{d^k}{dt^k} : B_r \rightarrow B_r$ acts on $\{u_j\}_{j \in \mathbb{Z}}$ by shift,

$$D^k u_j = u_{k+j}, \quad k > 0.$$

We write now $u_j = \sum_{n \geq 0} h_{n+j} \frac{t^n}{n!} = u_j(x_1, \dots, x_r; t)$, $j \in \mathbb{Z}$.

Operator $\partial_i = \partial / \partial x_i : B_r \rightarrow B_r$ induces a \mathbb{Q} -linear map

$$\partial_i : K_r \rightarrow K_r \quad (i = 1, \dots, r).$$

Operator $D^k = \frac{d^k}{dt^k} : B_r \rightarrow B_r$ acts on $\{u_j\}_{j \in \mathbb{Z}}$ by shift,

$$D^k u_j = u_{k+j}, \quad k > 0.$$

They both can be naturally extended first to operators in $\wedge^r K_r$, and then in F_0^r .

We denote them $\hat{\partial}_i$ and \hat{D}^k , resp.

(recall: $K_r = \text{Span}_{B_r} \{u_0, u_1, \dots, u_{1-r}\} = \text{Span}_{\mathbb{Q}} \{u_j, j \geq -r+1\}$).

Theorem 7: For the operators $\hat{\partial}_i$, $\hat{D}^k : F_0^r \rightarrow F_0^r$ ($1 \leq i, k \leq r$), we have

Theorem 7: For the operators $\hat{\partial}_i$, $\hat{D}^k : F_0^r \rightarrow F_0^r$ ($1 \leq i, k \leq r$), we have

$$c_0^r \hat{\partial}_i (c_0^r)^{-1} = \partial / \partial x_i,$$

$$c_0^r \hat{D}^k (c_0^r)^{-1} = kx_k,$$

$$[\hat{\partial}_i, \hat{D}^k] = [\partial / \partial x_i, kx_k] = i\delta_{ik}.$$

Theorem 7: For the operators $\hat{\partial}_i$, $\hat{D}^k : F_0^r \rightarrow F_0^r$ ($1 \leq i, k \leq r$), we have

$$c_0^r \hat{\partial}_i (c_0^r)^{-1} = \partial / \partial x_i,$$

$$c_0^r \hat{D}^k (c_0^r)^{-1} = kx_k,$$

$$[\hat{\partial}_i, \hat{D}^k] = [\partial / \partial x_i, kx_k] = i\delta_{ik}.$$

Denote \mathcal{H}_r the Lie algebra over \mathbb{Q} generated by $\{p_i, \hbar\}_{-r \leq i \leq r}$ such that

$$[p_i, p_k] = i\delta_{i,-k}\hbar, \quad [p_i, \hbar] = 0. \quad \dim \mathcal{H}_r = 2r + 2.$$

Theorem 7: For the operators $\hat{\partial}_i$, $\hat{D}^k : F_0^r \rightarrow F_0^r$ ($1 \leq i, k \leq r$), we have

$$c_0^r \hat{\partial}_i (c_0^r)^{-1} = \partial / \partial x_i,$$

$$c_0^r \hat{D}^k (c_0^r)^{-1} = kx_k,$$

$$[\hat{\partial}_i, \hat{D}^k] = [\partial / \partial x_i, kx_k] = i\delta_{ik}.$$

Denote \mathcal{H}_r the Lie algebra over \mathbb{Q} generated by $\{p_i, \hbar\}_{-r \leq i \leq r}$ such that

$$[p_i, p_k] = i\delta_{i,-k}\hbar, \quad [p_i, \hbar] = 0. \quad \dim \mathcal{H}_r = 2r + 2.$$

\mathcal{H}_∞ is the *oscillator Heisenberg Algebra* generated over \mathbb{C} by $\{p_i, \hbar\}_{i \in \mathbb{Z}}$ such that

$$[p_i, p_k] = i\delta_{i,-k}\hbar, \quad [p_i, \hbar] = 0.$$

We obtain two one-parametric families of representations of the Lie algebra \mathcal{H}_r :

We obtain two one-parametric families of representations of the Lie algebra \mathcal{H}_r :

Bosonic representations $\beta_m : B_r \rightarrow B_r$,

$$\forall P \in B_r : \quad \beta_m(p_0)P = mP, \quad \beta_m(\hbar)P = \hbar P, \quad \beta_m(p_j)P = \begin{cases} \partial P / \partial x_j, & j > 0 \\ -\hbar j x_{-j}, & j < 0 \end{cases}.$$

We obtain two one-parametric families of representations of the Lie algebra \mathcal{H}_r :

Bosonic representations $\beta_m : B_r \rightarrow B_r$,

$$\forall P \in B_r : \quad \beta_m(p_0)P = mP, \quad \beta_m(\hbar)P = \hbar P, \quad \beta_m(p_j)P = \begin{cases} \partial P / \partial x_j, & j > 0 \\ -\hbar j x_{-j}, & j < 0 \end{cases}.$$

Fermionic representations $\varphi_m : F_0^r \rightarrow F_0^r$,

$$\forall \Phi_\lambda^r = \bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots \in F_0^r :$$

$$\varphi_m(p_0)\Phi_\lambda^r = m \Phi_\lambda^r, \quad \varphi_m(\hbar)\Phi_\lambda^r = \hbar \Phi_\lambda^r, \quad \varphi_m(p_j)\Phi_\lambda^r = \begin{cases} \hat{\partial}_j \Phi_\lambda^r, & j > 0 \\ \hbar \hat{D}^{-j} \Phi_\lambda^r, & j < 0 \end{cases}.$$

We obtain two one-parametric families of representations of the Lie algebra \mathcal{H}_r :

Bosonic representations $\beta_m : B_r \rightarrow B_r$,

$$\forall P \in B_r : \quad \beta_m(p_0)P = mP, \quad \beta_m(\hbar)P = \hbar P, \quad \beta_m(p_j)P = \begin{cases} \partial P / \partial x_j, & j > 0 \\ -\hbar j x_{-j}, & j < 0 \end{cases}.$$

Fermionic representations $\varphi_m : F_0^r \rightarrow F_0^r$,

$$\forall \Phi_\lambda^r = \bar{u}_\lambda \wedge u_{-r} \wedge u_{-r-1} \wedge \dots \in F_0^r :$$

$$\varphi_m(p_0)\Phi_\lambda^r = m \Phi_\lambda^r, \quad \varphi_m(\hbar)\Phi_\lambda^r = \hbar \Phi_\lambda^r, \quad \varphi_m(p_j)\Phi_\lambda^r = \begin{cases} \hat{\partial}_j \Phi_\lambda^r, & j > 0 \\ \hbar \hat{D}^{-j} \Phi_\lambda^r, & j < 0 \end{cases}.$$

On this way, finite-dimensional counterparts of vertex operators appear
 [L. Gatto and P. Salehyan, arXiv 13.10.5132]